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Theoretical Interpretation of J-PARC E27 Data

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Study of kaonic nuclei by the d (π^+ , K⁺) reaction at J-PARC

Ichikawa al, XV International Conference on Hadron Spectroscopy-Hadron 2013 Nara, Japan

Simplest kaonic nucleus K⁻ pp bound state is searched at

J-PARC K 1.8 beam line (J-PARC E27 experiment)

d (π^+ , K⁺) X at the beam momentum 1.7 GeV/c X strangeness (-1), singly charged, double baryon K⁻ pp bound stae quasi free Λ^* p, Σ^* p, Λ p, Σ^0 p Anti Kaons \overline{K}^0 , $K^ \overline{d} s$, $\overline{u} s$ Strong attraction of $\overline{K} N$ in I = 0 channel

 Λ^* or Λ (1405)

A resonace lies below the K⁻ p threshold and in the continuum region of $\pi^- \Sigma^+$



"**∆(1405) Ansatz**"



DAΦNE Conf. (1999) Y. Akaishi & T. Yamazaki, Phys. Rev. C <u>65</u> (2002) 044005 T. Yamazaki & Y. Akaishi, Phys. Lett. B <u>535</u> (2002) 70



(В Е, Г) 115⁺⁶₋₅, 67⁺¹⁴₋₁₁ MeV

 $p + p \rightarrow \Lambda + p + K^{+}$ $T_{p} = 2.85 \text{ GeV}$ $103^{+3}_{-3}, 118^{+8}_{-8} \text{ MeV}$



X \longrightarrow strangeness (-1), single charged K⁻ pp, Λ^* p, Σ^* p, Σ^0 p

D (π⁺ , K⁺) X

Three-Body Coupled Channel Rearrangement Gaussian Basis Treatment



$$\psi(r,\rho) = \sum_{c=1}^{3} \sum_{i,j} \sum_{l,L} A_{c}^{i,j} r_{c}^{l+1} e^{-\left(\frac{r_{c}}{b_{i}}\right)^{2}} \rho_{c}^{L+1} e^{-\left(\frac{\rho_{c}}{d_{j}}\right)^{2}}$$

(1) and (2) are $\Lambda^* - p$ structure, strong KN interaction. Major contribution

(3) is K-(pp) structure with p-p repulsive interaction

$$H \approx H_1 = \left\{ T(\vec{r}) + V_{K^- p}(\vec{r}) \right\} + \left\{ T(\vec{\rho}) + V_{\Lambda^* p}(\vec{\rho}) \right\}$$
$$H_{\Lambda^*} + H_{\Lambda^* p}$$

Variational wave function of K-pp

ATMS

Amalgamation of Two-body correlations into Multiple Scattering process

$$\psi = [\phi_{12} + \phi_{13}]T = \frac{1}{2}$$

 $\Psi = \left[\left\{ f^{\prime=0}(r_{12}) \hat{P}_{12}^{\prime=0} + f^{\prime=1}(r_{12}) \hat{P}_{12}^{\prime=1} \right\} f_{NN}(r_{23}) f(r_{31}) + f(r_{12}) f_{NN}(r_{23}) \left\{ \frac{f^{\prime=0}(r_{31}) \hat{P}_{31}^{\prime=0}}{f^{\prime=0}(r_{31}) \hat{P}_{31}^{\prime=0}} + f^{\prime=1}(r_{31}) \hat{P}_{31}^{\prime=1} \right\} \right] |T = 1/2 \rangle$

$$\hat{P}_{12}^{I=0} = \frac{1 - \vec{\tau}_{\rm K} \vec{\tau}_{\rm N}}{4}, \quad \hat{P}_{12}^{I=1} = \frac{3 + \vec{\tau}_{\rm K} \vec{\tau}_{\rm N}}{4}$$

$$|T = 1/2\rangle = \sqrt{\frac{3}{4}} \left[\left(\overline{K}_1 N_2 \right)^{0,0} p_3 \right] + \sqrt{\frac{1}{4}} \left[-\sqrt{\frac{1}{3}} \left(\overline{K}_1 N_2 \right)^{1,0} p_3 + \sqrt{\frac{2}{3}} \left(\overline{K}_1 N_2 \right)^{1,1} n_3 \right]$$

$$\Lambda^* p$$

Euler-Lagrange equation

 $\delta_{f}\left\{\left\langle \Psi|\boldsymbol{H}|\boldsymbol{\Psi}\right\rangle-\lambda\left\langle \Psi|\boldsymbol{\Psi}\right\rangle\right\}=0$

$$v_{\rm KN}^{T=0}(r) = \{-595 - i83\}_{\rm MeV} \exp\{-(r/0.66_{\rm fm})^2\}$$
$$v_{\rm KN}^{T=1}(r) = \{-175 - i105\}_{\rm MeV} \exp\{-(r/0.66_{\rm fm})^2\}$$
$$v_{\rm NN}(r) = 2000_{\rm MeV} \exp\{-(r/0.447_{\rm fm})^2\} - 270_{\rm MeV} \exp\{-(r/0.942_{\rm fm})^2\} - 5_{\rm MeV} \exp\{-(r/2.5_{\rm fm})^2\}$$

Heitler-London-Heisenberg picture of K-pp



Model: K⁻ pp production as a $\Lambda(1405)$ doorway state

Y. Akaishi, T. Yamazaki, M. Obu and M. Wada, Nucl. Phys. A 835 (2010) 67

Molecular

Heitler-London (1927) Heisenberg (1932)

Nuclear Force

Yukawa (1935)

р

n

Super Strong **Nuclear Force** (2007)



Interaction between Λ^* and proton

 $V_{\Lambda^* p}$ is constructed from three-body ATMS calculation

Amalgamation of Two-body correlations into Multiple Scattering process

DISTO

 $B_{K^{-}pp} = 105 \text{ MeV}, \quad \Gamma = 118 \text{ MeV}$ $B_{\Lambda^{*}p} = 105 - 27 = 78 \text{ MeV}$ $V_{\Lambda^{*}p} = (V_{0} + iW_{0}) \left(\frac{r}{b}\right)^{2} e^{-\frac{r}{b}} \qquad b = 0.3 \text{ fm}$ $V_{0} = -400 \text{ MeV}$ $W_{0} = -162 \text{ MeV}$

L= 0 E = - 78.0 - i 59.3 MeV

On the missing mass spectrum from $D(\pi^+, K^+)$ E27 experiment



$$M_X^2 = (E_i - E_K)^2 - (p_\pi - p_K)^2$$
$$E_i = E_{\pi^+} + M_d$$

Differential cross section contains kinematical factor and spectral function S(E) where E is the missing mass variable

 $Y(\Lambda / \Sigma) \quad Y^*(\Lambda^* / \Sigma^*) \qquad S(E) = -\frac{1}{\pi} \operatorname{Im} \left[\int dr \, dr' f^*(r') G^{(+)}(r',r) f(r) \right]$

$$G^{(+)}(r',r) = \left\langle r' \left| \frac{1}{E - H_{K^-pp} + i\varepsilon} \right| r \right\rangle \qquad f(r) = e^{iQr} \psi_i(r)$$

$$E - H_{K^-pp} = (E - H_{\Lambda^*}) - H_{\Lambda^*p}$$

$$(E - H_{\Lambda^*}) \rightarrow E - B_{\Lambda^*} + i \frac{1}{2} \Gamma_{\Lambda^*}$$

$$G^{(+)}(r',r) = \left\langle r' \left| \frac{1}{E' - H_{\Lambda^*p} + i\varepsilon} \right| r \right\rangle$$

<u>Missing mass spectrum of Λ^* -p system</u>

for E27@J-PARC



Coincidence study









FINUDA

M. Agnello et al., Phys. Rev. Lett. <u>94</u> (2005) 212303



QF Λ -p production and Fermi motion

Deuteron momentum distribution $\rho_d(k_n) = N \times e^{-\frac{2}{a}k_n^2}$ where a = 0.1994 fm⁻¹

Neutron distribution is largest at $p_n = 0$ $\pi^+ + n \rightarrow K^+ + \Lambda$



$$M_X^2 = (E_{\pi} + M_d - E_K)^2 - (p_{\pi} - p_K)^2$$

Peak position = 2080.94 MeV, level width = 35 MeV consistent with exp.

 Λ peak is not a dynamical resonance

Quasi-free $\Lambda^* p$ production



Peak position 2440 MeV, 40 MeV shift from E27 exp.

Blue curve: Fermi motion suppressed, $\Gamma_{\Lambda}^{*} = 50 \text{ MeV}$ $\Gamma_{width} = 100 \text{ MeV}$ Pink curve: Fermi motion included, $\Gamma_{\Lambda}^{*} = 0$ $\Gamma_{fermi} = 90 \text{ MeV}$ Green curve: Fermi motion included, $\Gamma_{\Lambda}^{*} = 50 \text{ MeV}$ $\Gamma_{total} = 140 \text{ MeV}$

Concluding remarks

The $\Lambda^*=\Lambda(1405)$ plays an essential role in forming "anti-Kaonic Nuclear Clusters", the simplest one of which is K-pp = (K-p)-p = Λ^* -p.

The Λ^* -p structure interacting with "super-strong force" due to K^{bar} migration provides a possible explanation of recent J-PARC data on K⁻pp.



Thank you very much!